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Grain size effects on oxygen permeation in submicrometric $CaTi_{0.8}Fe_{0.2}O_{3−δ}$ ceramics obtained by mechanical activation

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Abstract

Ceramic samples of CaTi_{0.8}Fe_{0.2}O_{3−δ} were obtained from mechanically activated mixtures of TiO₂, Fe₂O₃ and CaCO₃ sintered at 1150 °C for 2 h. The ceramics are dense with submicrometric grains with sizes in the order of 100–200 nm and 400–600 nm. X-ray diffraction analysis revealed a single perovskite phase structure indicating a high level of homogeneity, which was confirmed by transmission electron microscopy. The oxygen permeability through these samples, measured in the range 700–950 ℃, is about 50% lower than for ceramics with considerably larger grains, close to 10 μ m. The permeability increases linearly with the increase of the grain size in the range from ≈200 nm to ≈10 μ m. Observations by results suggest that the grain boundaries have a negative impact on the high temperature ionic transport properties of CaTi_{0.8}Fe_{0.2}O_{3−δ} ceramics.

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1. Introduction

Iron substituted calcium titanate is a candidate for application as a ceramic membrane for oxygen separation due to an attractive combination of properties, including ionic and electronic conductivities, thermal expansion, thermo and chemi-cal stability and low cost.^{[1–6](#page-3-0)} The substitution of Ti^{4+} by Fe up to approximately 20 mol.% is compensated by oxygen vacancies; on the other hand, electrons are injected at low oxygen partial pressure, and electron holes (coupled to the existence of Fe^{4+} in oxidising conditions) in both cases leading to increased electronic conductivity.^{1–8} The oxygen vacancies tend to order around tetracoordinated Fe³⁺ for $x \ge \infty 0.2^{9-11}$ and a marked decrease in the oxygen conductivity, and thus in oxygen permeability, is observed.^{[5,6,8,10](#page-3-0)}

The material is usually obtained by conventional ceramic route which have a tendency to yield agglomerated powders with compositional inhomogeneities and demands for relatively high sintering temperatures (\geq 1320 °C) in order to obtain an impervious ceramic body.[1,2,7,12](#page-3-0) Moreover, coreshell grains tend to develop in CaTi_{0.8}Fe_{0.2}O_{3−δ}, with grain interiors consisting of pure $CaTiO₃$ and the periphery of an iron-substituted phase, representing an additional uncontrolled factor.¹²

On the other hand, the interest for nanosized ceramics has been increasing as they may present properties different from those of normal ceramics with grain sizes in the micron range. The mechanochemical synthesis is certainly an attractive alternative to obtain nano-, or submicrometric, ceramics. Pure CaTiO₃ was obtained by grinding TiO₂ and CaO or CaCO₃.^{[13](#page-3-0)} Alternatively, the simple mechanical activation may suffice to bring the precursor initial thermodynamic state to a far from equilibrium state, usually amorphous, which can then be heat treated to obtain the desired phase at a temperature lower than that needed in a conventional ceramic route.

This article reports on the synthesis of CaTi_{0.8}Fe_{0.2}O_{3−δ} from mechanically activated $TiO₂$, Fe₂O₃ and CaCO₃ precursors aiming to obtain dense, homogeneous CaTi $_{0.8}Fe_{0.2}O_{3-\delta}$ submicrometric ceramics and on subsequent oxygen perme-

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ation studies. Results are presented in comparison to those obtained for ceramics with larger grains obtained via the conventional ceramic route.[5,8](#page-3-0)

2. Experimental

The appropriate amounts of high purity powders of anatase TiO₂, CaCO₃ (Merck) and Fe₂O₃ (Riedel-de-Haën) were mixed in ethanol using a ball mill during 2 h at 20 rpm and dried in air at 60° C. These mixtures were then mechanically activated at room temperature by dry grinding in a Philips PW 4018 planetary ball mill at 500 rpm. The reactants were milled in a 45 cm^3 tetragonal stabilised zirconia (TZP) container using TZP balls (8 mm in diameter) with a ball to powder weight ratio of 10:1. The grinding was interrupted every 20 min to remove the powder from the wall of the container. Ceramic samples were obtained from powders isostatically pressed at 200 MPa and subsequent sintering in air at a maximum temperature of 1150° C for two hours with two additional temperature dwells of 1 h at 300 and 700 \degree C; the heating rate was 2 K/min. The density, measured by immersion in Hg, was higher than 93% of the theoretical density. The microstructure was analysed by SEM scanning electron microscopy (Hitachi S4100). The structure and phase purity was monitored in situ at various temperatures by powder X-ray diffraction (X'Pert MPD Philips diffractometer with Cu $K\alpha$ X-radiation). A detailed description of the processing procedures is presented elsewhere.[14](#page-3-0)

The oxygen permeability was measured with the usual system consisting of an oxygen sensor and an oxygen electrochemical pump, both of stabilised zirconia, sealed onto the membrane under analysis.^{[4–6](#page-3-0)} The oxygen gradient across the sample membrane is established between atmospheric air $(Po₂ = 21 kPa)$, at the feed side $(P₂)$, and a lower value (typically from 15 down to 1.3 kPa), at the permeate side (P_1) , obtained by pumping oxygen out of the chamber. The zirconia sensor is used to monitor P_1 according to the Nernst law. In a steady state condition, the current in the zirconia pump (*I*pump) equals the ionic current entering the chamber through the sample and, thus, the molecular oxygen flux, $j_{02} = I_{\text{pump}}/(4FS)$, where *S* is the surface area of the membrane available for transport and *F* the faraday constant. The measurements were carried out at temperatures in the range 700–950 ◦C. The gas tightness of the ceramic samples was verified by forcing a flow of compressed air $(\approx 250 \text{ kPa})$ through the membrane under water and checking that no bubbles form on the permeate side. The existence of leaks through the seals was assessed by examining the transient response of the cell, as described in refs.[14–16](#page-3-0)

3. Results

Fig. 1 shows a powder XRD pattern of the mechanically activated precursor mixture collected in situ at 1000 ◦C. It can

Fig. 1. Powder X-ray diffraction pattern collected in situ at $1000\,^{\circ}$ C (the * peaks correspond to the platinum).

be seen that the perovskite phase is formed and no secondary phases are apparent. Note that for non-activated precursors the desired single perovskite material can only be obtained at temperatures higher than $1250\,^{\circ}\text{C}$ and over considerably large periods of time $(>20 h)$.^{1,4,10}

However, the density of the activated ceramics sintered at $1000\degree$ C is rather low. Slightly higher temperatures were thus used to avoid oxygen leaks through the membrane during the permeability measurements. The SEM microstructures presented in Fig. 2 show that, indeed, the density of ceramics obtained at 1150° C for 2 h is rather high and that the few existing pores are not percolated. This result indicates that the activated powders are thus clearly more reactive and represents a decrease in the densification temperature, when compared to the conventionally prepared samples, of about

Fig. 2. Microstructure of ceramics sintered at 1150 ◦C for 2 h.

 $200\degree C$ ^{[1,4-7,12](#page-3-0)} Moreover, it shows that at such low sintering temperature, the ceramics retain a grain size in the submicrometric range. A more detailed analysis revealed that the size of the majority of the grains falls in the 150–200 nm range, while some larger grains (about 500 nm) are also observed, probably as a result of agglomerates in the initial powders.

It may be concluded that the mechanical activation of the precursor is indeed suitable to obtain ceramics with improved sinterability while retaining the possibility of designing the microstructure by appropriate manipulation of the sintering conditions and/or the activation procedures.

Fig. 3 shows oxygen permeability data obtained at different temperatures for the 1mm thick membrane, which microstructure is shown in [Fig. 2. T](#page-1-0)he flux of molecular oxygen, measured between air and 2.1 kPa of oxygen at the permeate side, is \approx 7.33 × 10⁻⁹ mol⁻¹ cm⁻² at 900 °C and about one order of magnitude lower at 700 ◦C. These values are considerably lower than those obtained for ceramics with larger grains. For example, the flux is about 30% higher at 900 °C for ceramics with an average grain size of 9 μ m.^{[1–6,12](#page-3-0)} Such trend is observed in a relatively broad temperature range, and follow an Arrhenius-like behaviour with apparent activation energy of 117 kJ mol−¹ (Fig. 4). This value is closer to that found for the ionic conduction in coarsegrained (\approx 10 µm) CaTi_{0.8}Fe_{0.2}O_{3−δ} and¹², thus, the oxygen permeability through the CaTi_{0.8}Fe_{0.2}O_{3−δ} submicrometric membranes should be mainly determined by the oxygen ion diffusion.

A linear relation is revealed when plotting the flux, for a given oxygen gradient, as a function of the average grain size of the ceramics (Fig. 5). Moreover, the slope increases with increasing temperature.

According to the brick-layer model (the simplest) for the conductivity of a polycrystalline material, the apparent conductance of the grain boundaries σ_{gb}^{app} is given by

$$
\sigma_{\text{gb}}^{\text{app}} = \sigma_{\text{gb}} \frac{G_{\text{av}}}{\delta_{\text{gb}}} \tag{1}
$$

Fig. 3. Oxygen electrochemical permeability measured at different temperatures.

Fig. 4. Temperature dependence of the oxygen permeability flux measured for ceramics with different grain size (results for the samples with larger G_{av} were taken from ref.^{[12](#page-3-0)}

where σ_{gb} and δ_{gb} are, respectively, the effective conductivity and thickness of an individual grain boundary and *G*av is the average equivalent grain diameter. On the other hand, if the reactions at the surface of the membrane are fast enough, jo2 is determined by the ambipolar transport properties of the material (σ _{amb}) and the external temperature (*T*) and Po₂ conditions:

$$
j_{02} = \frac{RT}{16F^2L} \int_{P_1}^{P_2} \sigma_{amb} \, d \ln P_{02}.
$$
 (2)

where *R* is the gas constant and *L* the membrane thickness. It follows from the combination of Eqs. (1) and (2), and under the assumption of grain boundary controlled transport, that the ratio of the ambipolar grain boundary conductivity and grain boundary thickness ($\sigma_{amb,gb}/\delta_{gb}$) is proportional to the slope (β) of the linear plots in Fig. 5:

$$
\frac{\sigma_{\rm amb,gb}}{\delta_{\rm gb}} = \frac{16\beta F^2}{RT \ln(P_2/P_1)}
$$
(3)

[Fig. 6](#page-3-0) shows estimates of $\sigma_{\text{amb.gb}}/\delta_{\text{gb}}$ (for $P_2/P_1 \approx 10$) in Arrhenius coordinates with an apparent activation energy, ca. 97 kJ/mol, which is similar to the value found for the total ionic conductivity measured for CaTi $_0$ ₈Fe $_0$ 2O_{3−δ} ceramics.[1,8,12](#page-3-0) Taking a further step on the interpretation

Fig. 5. Oxygen permeability flux as function of the average grain size *G*av (results for the samples with larger G_{av} were taken from ref.¹

Fig. 6. Temperature dependence of the ratio $\sigma_{amb,gb}/\delta_{gb}$ estimated for $P_2/P_1 \approx 10$.

of the results, one may consider previously obtained estimates of δ_{gb} , ca. 20 nm¹⁷ to obtain values for $\sigma_{amb,gb}$ of 8.1 × 10⁻⁶ S/cm at 850 °C or 1.7 × 10⁻⁵ S/cm at 950 °C. These values, although affected by the uncertainty in $\delta_{\alpha b}$, seem reasonable when compared to values, extrapolated to the same temperatures from impedance spectroscopy data, for Fe-doped SrTiO₃ (2.5 × 10⁻¹⁰ and 6.9 × 10⁻¹⁰ S/cm)¹⁸ and Ga-substituted CeO₂ (2.0 × 10⁻³ and 4.4 × 10⁻³ S/cm).¹⁹ Therefore, and assuming that σ_{amb} is determined mainly by the ionic conductivity^{1,12} and that the conductivity of the grain interior is not influenced by the grain size, the results now reported show that the oxygen transport in CaTi_{0.8}Fe_{0.2}O_{3−δ} is considerably constrained at the grain boundaries.

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